

ASYMPTOTIC BEHAVIOR FOR A STRUCTURE WITH INTERFACIAL SLIP AND ONE DISCONTINUOUS LOCAL INTERNAL KELVIN-VOIGT DAMPING

OCTAVIO P. VERA VILLAGRÁN AND CARLOS A. RAPOSO

ABSTRACT. In the present paper, we study the stabilization of laminated beams with one discontinuous local internal viscoelastic damping of Kelvin-Voigt type. First, we prove the strong stability of the system using the Arendt-Batty Theorem. Finally, we present the polynomial stability by using Borichev-Tomilov's result of optimality.

1. INTRODUCTION

In this work we study the stability for a structure with interfacial slip and frictional damping, given by,

$$\begin{cases} \rho_1 u_{tt} + k(\psi - u_x)_x - k d(x)(s_{xt} - u_t) = 0, \\ \rho_2 (s - \psi)_{tt} - b(s - \psi)_{xx} - k(\psi - u_x) + \beta (s - \psi)_t = 0, \\ \rho_2 s_{tt} - b s_{xx} + 3k(\psi - u_x) - 3k[d(x)(s_{xt} - u_t)]_x + 4\delta s_t = 0, \end{cases} \quad (1.1)$$

where $(x, t) \in (0, L) \times (0, +\infty)$. The coefficients $\rho_1, \rho_2, k, b, \gamma$ and β are positive and denote the density of the beam, the mass moment of inertia, the shear of stiffness, the flexural rigidity, the effective damping of the rotational angle and the adhesive damping parameter respectively. The function $u = u(x, t)$ denotes the transverse displacement, $\psi = \psi(x, t)$ represents the rotation angle, and $s = s(x, t)$ is proportional to the amount of slip along the interface at time t and longitudinal spatial variable x respectively.

We consider the following boundary conditions:

$$u(0, L) = u(L, t) = \psi(0, t) = \psi(L, t) = s(0, t) = s(L, t) = 0 \quad (1.2)$$

and initial data

$$\begin{cases} (u(x, 0), \psi(x, 0), s(x, 0)) = (u_0(x), \psi_0(x), s_0(x)), \\ (u_t(x, 0), \psi_t(x, 0), s_t(x, 0)) = (u_1(x), \psi_1(x), s_1(x)). \end{cases} \quad (1.3)$$

2010 *Mathematics Subject Classification.* 35Q53, 35Q55, 47J353, 35B35.

Key words and phrases. Polynomial stability, strong stability, Kelvin-Voigt damping.

Submitted Feb. 24, 2021.

We suppose that there exists $0 < \alpha < \beta < L$ and a positive constant d_0 such that

$$d(x) = \begin{cases} d_0 & \text{if } x \in (\alpha, \beta), \\ 0 & \text{if } x \in (0, \alpha) \cup (\beta, L). \end{cases} \quad (1.4)$$

The model (1.1) is closely related with the Timoshenko's theory. Timoshenko's theory started in 1921 [25, 26] and since then, has been extensively studied by several authors over different points of view. Most of them studied the global well-posedness, asymptotic behavior among other properties (see [1, 4, 9, 10, 11] and references therein).

Hansen in 1994 (see [12, 13] for model description) derived from Timoshenko's theory the model (1.1), that describes a structure of two identical beams of uniform thickness with an adhesive layer (of negligible thickness and mass) bonding the two adjoining surfaces in such a way that a slip is allowed while they are continuously in contact with each other. The model (1.1) is called laminated beams and has gained a lot of interest in recent years. We mention for instance [2, 6, 7, 8, 15, 16, 17, 19, 20, 22, 23, 24] and references therein.

In this work, we have obtained the polynomial stability for a structure with interfacial slip with one discontinuous local internal Kelvin-Voigt damping. In order to obtain the results we will use the theorem given by Arendt-Batty [3] together with a result given by Borichev-Tomilov [5]. Indeed, our main results is:

Theorem 1.1. *There is a constant $C > 0$ such that for every $U_0 \in \mathcal{D}(\mathcal{A})$, we have*

$$E(t) \leq \frac{C}{t} \|W_0\|_{\mathcal{D}(\mathcal{A})}, \quad t > 0. \quad (1.5)$$

This paper is organized as follows. In Section 2 we recall some auxiliary results. Section 3 by semigroup theory of linear operators we obtain an existence theorem of solutions of system (2.7)-(2.9). In Section 4, we show the strong stability by using Arendt-Batty's theorem. In Section 5, we present the polynomial stability of the corresponding semigroup using Borichev-Tomilov's result of optimality.

2. PRELIMINARY

Throughout this paper, C is a generic constant, not necessarily the same at each occasion (it will change from line to line) and depends on the indicated quantities. In the following lemma we prove the dissipative properties of the system (1.1) in the sense that its energy is non-decreasing with respect to time, that is,

Lemma 2.1. *(Energy of the system) For every solution (u, ψ, s) of the system (1.1), the total energy $\mathcal{E} : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ defined at time t by*

$$\mathcal{E}(t) = \frac{1}{2} \left[3\rho_1 \|u_t\|_{L^2(0, L)}^2 + 3k \|\psi - u_x\|_{L^2(0, L)}^2 + \rho_2 \|s_t\|_{L^2(0, L)}^2 + b \|s_x\|_{L^2(0, L)}^2 + 3\rho_2 \|(s - \psi)_t\|_{L^2(0, L)}^2 + 3b \|(s - \psi)_x\|_{L^2(0, L)}^2 \right] \quad (2.1)$$

satisfies

$$\begin{aligned} \frac{d}{dt} \mathcal{E}(t) &= -3k \int_0^L d(x) |s_{xt} - u_t|^2 dx - 3\beta \int_0^L |(s - \psi)_t|^2 dx - 4\delta \int_0^L |s_t|^2 dx \\ &= -3k d_0 \int_\alpha^\beta |s_{xt} - u_t|^2 dx - 3\beta \int_0^L |(s - \psi)_t|^2 dx - 4\delta \int_0^L |s_t|^2 dx. \end{aligned} \quad (2.2)$$

Proof. Let $(u, u_t, \psi, \psi_t, s, s_t)$ be a regular solution of system (1.1). Multiplying (1.1)₁ by $3u_t$ we have

$$\frac{1}{2} \frac{d}{dt} 3\rho_1 \int_0^L u_t^2 dx + 3k \int_0^L (\psi - u_x)_x u_t dx - 3k \int_0^L d(x) (s_{xt} - u_t) u_t dx = 0.$$

Using integration by parts and (1.2) it follows that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} 3\rho_1 \int_0^L u_t^2 dx - 3k \int_0^L (\psi - u_x) u_{xt} dx \\ - 3k \int_0^L d(x) s_{xt} u_t dx + 3k \int_0^L d(x) u_t^2 dx = 0. \end{aligned}$$

The second term in the last equation can be written as

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} 3\rho_1 \int_0^L u_t^2 dx - 3k \int_0^L (\psi - u_x) (u_x - \psi + \psi)_t dx \\ - 3k \int_0^L d(x) s_{xt} u_t dx + 3k \int_0^L d(x) u_t^2 dx = 0. \end{aligned}$$

Then

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} 3\rho_1 \int_0^L u_t^2 dx + 3k \int_0^L (\psi - u_x) (\psi - u_x)_t dx \\ - 3k \int_0^L (\psi - u_x) \psi_t dx - 3k \int_0^L d(x) s_{xt} u_t dx + 3k \int_0^L d(x) u_t^2 dx = 0. \end{aligned}$$

Thus

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left[3\rho_1 \|u_t\|_{L^2(0,L)}^2 + 3k \|\psi - u_x\|_{L^2(0,L)}^2 \right] \\ - 3k \int_0^L (\psi - u_x) \psi_t dx - 3k \int_0^L d(x) s_{xt} u_t dx + 3k \int_0^L d(x) u_t^2 dx = 0. \end{aligned} \quad (2.3)$$

Now, multiplying (1.1)₃ by s_t and integrating over $(0, L)$ yields

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \rho_2 \int_0^L s_t^2 dx - b \int_0^L s_{xx} s_t dx \\ + 3k \int_0^L (\psi - u_x) s_t dx - 3k \int_0^L [d(x) (s_{xt} - u_t)]_x s_t + 4\delta \int_0^L |s_t|^2 dx = 0. \end{aligned}$$

Integrating by parts and using (1.2) it follows that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \rho_2 \int_0^L s_t^2 dx + b \int_0^L s_x s_{xt} dx \\ + 3k \int_0^L (\psi - u_x) s_t dx + 3k \int_0^L d(x) (s_{xt} - u_t) s_{xt} dx + 4\delta \int_0^L |s_t|^2 dx = 0. \end{aligned}$$

Then

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \left[\rho_2 \|s_t\|_{L^2(0,L)}^2 + b \|s_x\|_{L^2(0,L)}^2 \right] + 3k \int_0^L (\psi - u_x) s_t \, dx \\ & + 3k \int_0^L d(x) s_{xt}^2 \, dx - 3k \int_0^L d(x) u_t s_{xt} \, dx + 4\delta \int_0^L |s_t|^2 \, dx = 0. \end{aligned} \quad (2.4)$$

Finally, multiplying (1.1)₂ by $3(s - \psi)_t$ and integrating over $(0, L)$ we have

$$\begin{aligned} & 3\rho_2 \int_0^L (s - \psi)_{tt} (s - \psi)_t \, dx - 3b \int_0^L (s - \psi)_{xx} (s - \psi)_t \, dx \\ & - 3k \int_0^L (\psi - u_x) (s - \psi)_t \, dx + 3\beta \int_0^L |(s - \psi)_t|^2 \, dx = 0. \end{aligned}$$

Integrating by parts and using (1.2) it follows that

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} 3\rho_2 \int_0^L |(s - \psi)_t|^2 \, dx + \frac{1}{2} \frac{d}{dt} 3b \int_0^L |(s - \psi)_x|^2 \, dx \\ & - 3k \int_0^L (\psi - u_x) (s - \psi)_t \, dx + 3\beta \int_0^L |(s - \psi)_t|^2 \, dx = 0. \end{aligned}$$

Thus

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \left[3\rho_2 \|(s - \psi)_t\|_{L^2(0,L)}^2 + 3b \|(s - \psi)_x\|_{L^2(0,L)}^2 \right] \\ & - 3k \int_0^L (\psi - u_x) (s - \psi)_t \, dx + 3\beta \int_0^L |(s - \psi)_t|^2 \, dx = 0. \end{aligned} \quad (2.5)$$

Adding (2.3), (2.4) and (2.5) the Lemma 2.1 follows. \square

For convenience, from now and on, we introduce a new variable z given by

$$z = s - \psi \quad (\text{the effective rotation angle}) \iff \psi = s - z, \quad (2.6)$$

then (2.5) becomes

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \left[3\rho_2 \|z_t\|_{L^2(0,L)}^2 + 3b \|z_x\|_{L^2(0,L)}^2 \right] \\ & - 3k \int_0^L (s - z - u_x) z_t \, dx + 3\beta \int_0^L |z_t|^2 \, dx = 0. \end{aligned}$$

Performing the change of variable (2.6), system (1.1) leads to

$$\begin{cases} \rho_1 u_{tt} + k(s - z - u_x)_x - k d(x) (s_{xt} - u_t) = 0, \\ \rho_2 z_{tt} - b z_{xx} - k(s - z - u_x) + \beta z_t = 0, \\ \rho_2 s_{tt} - b s_{xx} + 3k(s - z - u_x) - 3k[d(x)(s_{xt} - u_t)]_x + 4\gamma s + 4\delta s_t = 0, \end{cases} \quad (2.7)$$

with the boundary conditions

$$u(0, L) = u(L, t) = z(0, t) = z(L, t) = s(0, t) = s(L, t) = 0 \quad (2.8)$$

and initial data

$$\begin{cases} (u(x, 0), z(x, 0), s(x, 0)) = (u_0(x), z_0(x), s_0(x)), \\ (u_t(x, 0), z_t(x, 0), s_t(x, 0)) = (u_1(x), z_1(x), s_1(x)). \end{cases} \quad (2.9)$$

We present the following definition.

Definition 2.2. [18] Assume that \mathcal{A} is the generator of C_0 -semigroup of contractions $(e^{t\mathcal{A}})_{t \geq 0}$ on a Hilbert space \mathcal{H} . The C_0 -semigroup $(e^{t\mathcal{A}})_{t \geq 0}$ is said to be

a) Strongly stable if

$$\lim_{t \rightarrow +\infty} \|e^{t\mathcal{A}} x_0\|_{\mathcal{H}} = 0, \quad \forall x_0 \in \mathcal{H}. \quad (2.10)$$

b) Exponentially (or uniformly) stable if there exists two positive constant M and ε such that

$$\|e^{t\mathcal{A}} x_0\|_{\mathcal{H}} \leq M e^{-\varepsilon t} \|x_0\|_{\mathcal{H}}, \quad \forall t > 0, \quad \forall x_0 \in \mathcal{H}. \quad (2.11)$$

c) Polynomially stable if there exists two positive constant C and α such that

$$\|e^{t\mathcal{A}} x_0\|_{\mathcal{H}} \leq C t^{-\alpha} \|\mathcal{A}x_0\|_{\mathcal{H}}, \quad \forall t > 0, \quad \forall x_0 \in \mathcal{D}(\mathcal{A}). \quad (2.12)$$

The following results will use some time from now on and are fundamental to the proof of our theorems of stability.

Theorem 2.3. [3] Assume that \mathcal{A} is the generator of a C_0 -semigroup of contractions $(e^{t\mathcal{A}})_{t \geq 0}$ on a Hilbert space \mathcal{H} . If \mathcal{A} has no pure imaginary eigenvalues and $\sigma(\mathcal{A}) \cap i\mathbb{R}$ is countable, where $\sigma(\mathcal{A})$ denote the spectrum of \mathcal{A} , then the C_0 -semigroup $(e^{t\mathcal{A}})_{t \geq 0}$ is strongly stable.

Theorem 2.4. [5] Assume that \mathcal{A} is the generator of a strongly continuous semigroup of contractions $(e^{t\mathcal{A}})_{t \geq 0}$ on \mathcal{H} . If $i\mathbb{R} \subset \varrho(\mathcal{H})$, then for a fixed $\ell > 0$ the following conditions are equivalent

$$\sup_{\lambda \in \mathbb{R}} \left\| (i\lambda I - \mathcal{A})^{-1} \right\|_{\mathcal{L}(\mathcal{H})} = \mathcal{O}(|\lambda|^\ell),$$

$$\|e^{t\mathcal{A}} U_0\|_{\mathcal{H}}^2 \leq \frac{C}{t^{2/\ell}} \|U_0\|_{\mathcal{D}(\mathcal{A})}^2, \quad \forall t > 0, \quad U_0 \in \mathcal{D}(\mathcal{A}), \quad \text{for some } C > 0.$$

3. SETTING OF THE SEMIGROUP

In this section, we use results of the semigroup theory of linear operators to obtain an existence theorem of solutions of system (2.7)-(2.9) (see [21]). We introduce the Hilbert space

$$\mathcal{H} = [H_0^1(0, L) \times L^2(0, L)]^3,$$

equipped with the inner product $(u_t = U, \psi_t = \Psi, s_t = S$ with $z = s - \psi$ and $Z = S - \Psi)$

$$\begin{aligned} \langle W, \widetilde{W} \rangle_{\mathcal{H}} &= 3\rho_1 \int_0^L U \widetilde{U} \, dx + 3\rho_2 \int_0^L Z \widetilde{Z} \, dx + \rho_2 \int_0^L S \widetilde{S} \, dx \\ &+ b \int_0^L s_x \widetilde{s}_x \, dx + 3k \int_0^L (s - z - u_x) (\widetilde{s} - \widetilde{z} - \widetilde{u}_x) \, dx + 3b \int_0^L z_x \widetilde{z}_x \, dx, \end{aligned} \quad (3.1)$$

where $W = (u, u_t = U; z, z_t = Z; s, s_t = S) \in \mathcal{H}$, $\widetilde{W} = (\widetilde{u}, \widetilde{U}; \widetilde{z}, \widetilde{Z}; \widetilde{s}, \widetilde{S})$ and norm

$$\begin{aligned} \|W\|_{\mathcal{H}}^2 &= 3\rho_1 \|u_t\|_{L^2(0, L)}^2 + 3\rho_2 \|z_t\|_{L^2(0, L)}^2 + \rho_2 \|s_t\|_{L^2(0, L)}^2 \\ &+ b \|s_x\|_{L^2(0, L)}^2 + 3k \|s - z - u_x\|_{L^2(0, L)}^2 + 3b \|z_x\|_{L^2(0, L)}^2. \end{aligned} \quad (3.2)$$

Now, we rewrite the system (2.7) as an abstract problem on the Hilbert space \mathcal{H} given by

$$\frac{d}{dt}W(t) = \mathcal{A}W(t), \quad (3.3)$$

$$W(0) = W_0 = (u_0, u_1, z_0, z_1, s_0, s_1), \quad \forall t > 0. \quad (3.4)$$

The operator

$$\mathcal{A} : \mathcal{D}(\mathcal{A}) \subset \mathcal{H} \rightarrow \mathcal{H}$$

is given by

$$\mathcal{A} \begin{pmatrix} u \\ U \\ z \\ Z \\ s \\ S \end{pmatrix} = \begin{pmatrix} U \\ \frac{1}{\rho_1} [-k(s - z - u_x)_x + k d(x)(S_x - U)] \\ Z \\ \frac{1}{\rho_2} [b z_{xx} + k(s - z - u_x) - \beta Z] \\ S \\ \frac{1}{\rho_2} (b s_{xx} - 3k(s - z - u_x) + 3k[d(x)(S_x - U)]_x - 4\delta S) \end{pmatrix} \quad (3.5)$$

with domain

$$\mathcal{D}(\mathcal{A}) = \left\{ W \in \mathcal{H} : u, z, s \in H^2(0, L) \cap H_0^1(0, L), U, Z, S \in H_0^1(0, L), [d(x)(s_x - u)]_x \in L^2(0, L) \right\}.$$

To prove that \mathcal{A} is the infinitesimal generator of a C_0 -contraction semigroup, we consider the two following lemmas.

Lemma 3.1. *The operator \mathcal{A} is a dissipative.*

Proof.

$$\begin{aligned}
\langle \mathcal{A}W, W \rangle_{\mathcal{H}} &= 3\rho_1 \int_0^L \frac{1}{\rho_1} [-k(s-z-u_x)_x + k d(x)(S_x - U)] \bar{U} \, dx \\
&\quad + 3\rho_2 \int_0^L \frac{1}{\rho_2} [b z_{xx} + k(s-z-u_x) - \beta Z] \bar{Z} \, dx \\
&\quad + \rho_2 \int_0^L \frac{1}{\rho_2} [b s_{xx} - 3k(s-z-u_x) + 3k[d(x)(S_x - U)]_x - 4\delta S] \bar{S} \, dx \\
&\quad + 3k \int_0^L (S - Z - U_x)(\bar{s} - \bar{z} - \bar{u}_x) \, dx + 3k \int_0^L Z_x \bar{z}_x \, dx + b \int_0^L S_x \bar{s}_x \, dx \\
&= -3k \int_0^L (s-z-u_x)_x \bar{U} \, dx + 3k \int_0^L d(x) S_x \bar{U} \, dx - 3k \int_0^L d(x) |U|^2 \, dx \\
&\quad + 3b \int_0^L z_{xx} \bar{Z} \, dx + 3k \int_0^L (s-z-u_x) \bar{Z} \, dx - \beta \int_0^L |Z|^2 \, dx - 4\delta \int_0^L |S|^2 \, dx \\
&\quad + b \int_0^L s_{xx} \bar{S} \, dx - 3k \int_0^L (s-z-u_x) \bar{S} \, dx + 3k \int_0^L [d(x)(S_x - U)]_x \bar{S} \, dx \\
&\quad + b \int_0^L S_x \bar{s}_x \, dx + 3k \int_0^L (S - Z - U_x)(\bar{s} - \bar{z} - \bar{u}_x) \, dx + 3b \int_0^L Z_x \bar{z}_x \, dx.
\end{aligned}$$

Using (2.8) and straightforward calculations we have

$$\begin{aligned}
\langle \mathcal{A}W, W \rangle_{\mathcal{H}} &= 3k \int_0^L \left[(S - Z - U_x)(\bar{s} - \bar{z} - \bar{u}_x) - \overline{(S - Z - U_x)(\bar{s} - \bar{z} - \bar{u}_x)} \right] \, dx \\
&\quad + \int_0^L [Z_x \bar{z}_x - \overline{Z_x \bar{z}_x}] \, dx + b \int_0^L [S_x \bar{s}_x - \overline{S_x \bar{s}_x}] \, dx \\
&\quad - \beta \int_0^L |Z|^2 \, dx - 4\delta \int_0^L |S|^2 \, dx - 3k \int_0^L d(x) |S_x - U|^2 \, dx.
\end{aligned}$$

Then

$$\begin{aligned}
\langle \mathcal{A}W, W \rangle_{\mathcal{H}} &= 6ik \operatorname{Im} \int_0^L (S - Z - U_x)(\bar{s} - \bar{z} - \bar{u}_x) \, dx \\
&\quad + 2i \operatorname{Im} \int_0^L Z_x \bar{z}_x \, dx + 2ib \operatorname{Im} \int_0^L S_x \bar{s}_x \, dx \\
&\quad - \beta \int_0^L |Z|^2 \, dx - 4\delta \int_0^L |S|^2 \, dx - 3k \int_0^L d(x) |S_x - U|^2 \, dx.
\end{aligned}$$

Taking the real part we obtain

$$\begin{aligned}
\operatorname{Re} \langle \mathcal{A}W, W \rangle_{\mathcal{H}} &= -\beta \int_0^L |Z|^2 \, dx - 4\delta \int_0^L |S|^2 \, dx - 3k \int_0^L d(x) |S_x - U|^2 \, dx \\
&= -\beta \int_0^L |Z|^2 \, dx - 4\delta \int_0^L |S|^2 \, dx - 3k d_0 \int_{\alpha}^{\beta} |S_x - U|^2 \, dx \leq 0.
\end{aligned} \tag{3.6}$$

It follows that \mathcal{A} is a dissipative operator. \square

Lemma 3.2. *The operator \mathcal{A} is Maximal.*

Proof. Let $F = (f_1, f_2, f_3, f_4, f_5, f_6) \in \mathcal{H}$, we must show that there exists a unique $W = (u, u_t, z, z_t, s, s_t)^T \in \mathcal{D}(\mathcal{A})$ such that $\mathcal{A}W = F$, that is,

$$\begin{cases} u_t = f_1 & \text{in } H_0^1(0, L), \\ -k(s - z - u_x)_x + k d(x)(s_{xt} - u_t) = \rho_1 f_2 & \text{in } L^2(0, L), \\ z_t = f_3 & \text{in } H_0^1(0, L), \\ b z_{xx} + k(s - z - u_x) - \beta z_t = \rho_2 f_4 & \text{in } L^2(0, L), \\ s_t = f_5 & \text{in } H_0^1(0, L), \\ b s_{xx} - 3k(s - z - u_x) + 3k[d(x)(s_{xt} - u_t)]_x - 4\delta s_t = \rho_2 f_6 & \text{in } L^2(0, L), \end{cases} \quad (3.7)$$

with the boundary conditions

$$u(0, t) = u(L, t) = z(0, t) = z(L, t) = s(0, t) = s(L, t) = 0. \quad (3.8)$$

Replacing (3.7)_{1,3,5} into (3.7)_{2,4,6} respectively we have

$$\begin{cases} -k(s - z - u_x)_x = \rho_1 f_2 - k d(x)(f_{5x} - f_1) & \text{in } L^2(0, L), \\ b z_{xx} + k(s - z - u_x) = \rho_2 f_4 + \beta f_3 & \text{in } L^2(0, L), \\ b s_{xx} - 3k(s - z - u_x) + 3k[d(x)(f_{5x} - f_1)]_x = \rho_2 f_6 + 4\delta f_5 & \text{in } L^2(0, L). \end{cases} \quad (3.9)$$

Let $(\Phi_1, \Phi_2, \Phi_3) \in [H_0^1(0, L)]^3$. Then multiplying (3.9)_{1,2,3} by $\bar{\Phi}_1, \bar{\Phi}_2$ and $\bar{\Phi}_3$ respectively, integrating over $(0, L)$, using (3.8) and performing straightforward calculations we obtain

$$\begin{aligned} k \int_0^L (s - z - u_x) \bar{\Phi}_{1x} dx &= \rho_1 \int_0^L f_2 \bar{\Phi}_1 dx - k \int_0^L d(x)(f_{5x} - f_1) \bar{\Phi}_1 dx, \\ -b \int_0^L z_x \bar{\Phi}_{2x} dx + k \int_0^L (s - z - u_x) \bar{\Phi}_2 dx &= \rho_2 \int_0^L f_4 \bar{\Phi}_2 dx + \beta \int_0^L f_3 \bar{\Phi}_2 dx, \\ -b \int_0^L s_x \bar{\Phi}_{3x} dx - 3k \int_0^L (s - z - u_x) \bar{\Phi}_3 dx \\ &= \rho_2 \int_0^L f_6 \bar{\Phi}_3 dx - 3k \int_0^L d(x)(f_{5x} - f_1) \bar{\Phi}_{3x} dx + 4\delta \int_0^L f_5 \bar{\Phi}_3 dx. \end{aligned}$$

We define

$$\mathcal{B}((u, z, s), (\Phi_1, \Phi_2, \Phi_3)) = \mathcal{L}((\bar{\Phi}_1, \bar{\Phi}_2, \bar{\Phi}_3)) \quad (3.10)$$

for all $(\Phi_1, \Phi_2, \Phi_3) \in [H_0^1(0, L)]^3$, where

$$\begin{aligned} \mathcal{B}((u, z, s), (\Phi_1, \Phi_2, \Phi_3)) &= k \int_0^L (s - z - u_x) \bar{\Phi}_{1x} dx \\ &\quad - b \int_0^L z_x \bar{\Phi}_{2x} dx + k \int_0^L (s - z - u_x) \bar{\Phi}_2 dx \\ &\quad - b \int_0^L s_x \bar{\Phi}_{3x} dx - 3k \int_0^L (s - z - u_x) \bar{\Phi}_3 dx \end{aligned}$$

and

$$\begin{aligned} \mathcal{L}((\Phi_1, \Phi_2, \Phi_3)) &= \rho_1 \int_0^L f_2 \bar{\Phi}_1 dx - k \int_0^L d(x) (f_{5x} - f_1) \bar{\Phi}_1 dx \\ &\quad + \rho_2 \int_0^L f_4 \bar{\Phi}_2 dx + \beta \int_0^L f_3 \bar{\Phi}_2 dx + \rho_2 \int_0^L f_6 \bar{\Phi}_3 dx \\ &\quad - 3k \int_0^L d(x) (f_{5x} - f_1) \bar{\Phi}_{3x} dx + 4\delta \int_0^L f_5 \bar{\Phi}_3 dx. \end{aligned}$$

It is not difficult to see \mathcal{B} is a sesquilinear and continuous form on $[H_0^1(0, L)]^3 \times [H_0^1(0, L)]^3$ and \mathcal{L} is a linear and continuous form on $[H_0^1(0, L)]^3$. Thereby, \mathcal{B} is a coercive form on $[H_0^1(0, L)]^3 \times [H_0^1(0, L)]^3$. Hence, it follows by the Lax-Milgram Theorem that (3.10) admits a unique solution $(u, z, s) \in [H_0^1(0, L)]^3$. Taking test-functions $(\Phi_1, \Phi_2, \Phi_3) \in [\mathcal{D}(0, L)]^3$ we have that (3.9) holds in the distributional sense, from which we deduce that $(u, z, s) \in [H^2(0, L) \cap H_0^1(0, L)]^3$, while $[d(x)(f_{5x} - f_1)]_x \in L^2(0, L)$. Thereby, $U = (u, f_1, z, f_3, s, f_5)^T \in \mathcal{D}(\mathcal{A})$ is the unique solution of $\mathcal{A}U = F$. Thus \mathcal{A} is an isomorphism and since $\rho(\mathcal{A})$ is open set of \mathbb{C} (see [14]: Theorem 6.7, Chapter III), we easily get $R(\lambda I - \mathcal{A}) = \mathcal{H}$ for a sufficiently small $\lambda > 0$. This and using the dissipativity of \mathcal{A} , imply that $\mathcal{D}(\mathcal{A})$ is dense in \mathcal{H} and that is m-dissipative in \mathcal{H} (see [21]: Theorems 4.5, 4.6), the lemma is proved. \square

Proposition 3.3. *The operator \mathcal{A} is the infinitesimal generator of a C_0 -semigroup of contractions $\{\mathcal{S}_{\mathcal{A}}(t)\}_{t \geq 0}$.*

Proof. By the lemma 3.1, the operator \mathcal{A} is a dissipative and from lemma 3.2 we have that \mathcal{A} is maximal. Then, according the well known Lumer-Phillips Theorem [21] we have that the operator \mathcal{A} generates a C_0 -semigroup of contractions $e^{t\mathcal{A}}$ in \mathcal{H} . \square

Thereby we have the following result, which gives the well-posedness of (3.3)-(3.4).

Theorem 3.4. *For all $W_0 \in \mathcal{H}$, the system (3.3)-(3.4) admits a unique weak solution*

$$W(t) = e^{t\mathcal{A}} W_0 \in C^0(\mathbb{R}^+, \mathcal{H}). \quad (3.11)$$

Moreover, if $W_0 \in \mathcal{D}(\mathcal{A})$, then $W(t)$ is the unique strong solution with the following regularity

$$W(t) \in C^0(\mathbb{R}^+, \mathcal{D}(\mathcal{A})) \cap C^1(\mathbb{R}^+, \mathcal{H}). \quad (3.12)$$

4. STRONG STABILITY

In this section, we prove the strong stability (Theorem 4.3) of the system (2.7). The idea is to use the Theorem 2.3 due to Arendt-Batty [3]. According to Theorem 2.3, we will prove that the operator \mathcal{A} has no pure imaginary eigenvalues and $\sigma(\mathcal{A}) \cap i\mathbb{R}$ is countable.

Lemma 4.1. *For all $\lambda \in \mathbb{R}$, $(i\lambda I - \mathcal{A})$ is injective, that is,*

$$\text{Ker}(i\lambda I - \mathcal{A}) = \{0\}, \quad \forall \lambda \in \mathbb{R}.$$

Proof. By the Lemma 3.2, $0 \in \varrho(\mathcal{A})$. Then we will show the result for $\lambda \in \mathbb{R}^*$. For this aim, suppose that there exists a real number $\lambda \neq 0$ and $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ such that

$$\mathcal{A}W = i\lambda W. \quad (4.1)$$

In fact, we have

$$\begin{cases} U = i\lambda u, \\ -k(s - z - u_x)_x + k d(x)(S_x - U) = i\lambda \rho_1 U, \\ Z = i\lambda z, \\ b z_{xx} + k(s - z - u_x) - \beta Z = i\lambda \rho_2 Z, \\ S = i\lambda s, \\ b s_{xx} - 3k(s - z - u_x) + 3k[d(x)(S_x - U)]_x - 4\delta S = i\lambda \rho_2 S. \end{cases} \quad (4.2)$$

From (4.1) we have

$$\langle i\lambda W, W \rangle_{\mathcal{H}} = \langle \mathcal{A}W, W \rangle_{\mathcal{H}} \iff i\lambda \|W\|_{\mathcal{H}}^2 = \langle \mathcal{A}W, W \rangle_{\mathcal{H}}. \quad (4.3)$$

Taking the real part and using (3.6) we obtain

$$\begin{aligned} 0 &= \operatorname{Re} [i\lambda \|W\|_{\mathcal{H}}^2] \\ &= -\beta \int_0^L |Z|^2 dx - 4\delta \int_0^L |S|^2 dx - 3k d_0 \int_{\alpha}^{\beta} |S_x - U|^2 dx \leq 0. \end{aligned} \quad (4.4)$$

Thereby we have

$$\begin{cases} Z = 0 & \text{in } (0, L), \\ S = 0 & \text{in } (0, L), \\ S_x - U = 0 & \text{in } (\alpha, \beta). \end{cases} \quad (4.5)$$

Using (4.2)_{3,5} into (4.5)_{1,2} and then fact that $\lambda \neq 0$ we have

$$\begin{cases} z = 0 & \text{in } (0, L), \\ s = 0 & \text{in } (0, L). \end{cases} \quad (4.6)$$

Replacing (4.6)₂ into (4.5)₃ we have

$$U = 0 \quad \text{in } (\alpha, \beta). \quad (4.7)$$

Now, replacing (4.2)_{1,5} into (4.5)₃ and using that $\lambda \neq 0$ we obtain

$$s_x - u = 0 \iff s_x = u \quad \text{in } (\alpha, \beta). \quad (4.8)$$

Therefore from (4.7) into (4.8) we obtain

$$u = 0 \quad \text{in } (\alpha, \beta). \quad (4.9)$$

Thereby $W = (0, 0, 0, 0, 0, 0)$ in (α, β) .

On the other hand, (4.2) can be written in $(0, \alpha) \cup (\beta, L)$, that is,

$$\begin{cases} \lambda^2 \rho_1 u - k(s - z - u_x)_x = 0 & \text{in } (0, \alpha) \cup (\beta, L), \\ \lambda^2 \rho_2 z + b z_{xx} + k(s - z - u_x) = 0 & \text{in } (0, \alpha) \cup (\beta, L), \\ \lambda^2 \rho_2 s + b s_{xx} - 3k(s - z - u_x) = 0 & \text{in } (0, \alpha) \cup (\beta, L). \end{cases} \quad (4.10)$$

Let $\mathcal{W} = (u, u_x, z, z_x, s, s_x)^T$ and we define

$$\mathcal{W}_x = \mathcal{B}_\lambda \mathcal{W} \quad (4.11)$$

where differentiating in x -variable, $x \in (0, \alpha)$ we have

$$\underbrace{\begin{pmatrix} u_x \\ u_{xx} \\ z_x \\ z_{xx} \\ s_x \\ s_{xx} \end{pmatrix}}_{\mathcal{W}_x} = \underbrace{\begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 \\ -\frac{\lambda^2 \rho_1}{k} & 0 & 0 & -1 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & \frac{k}{b} & -\frac{(\lambda^2 \rho_2 - k)}{b} & 0 & -\frac{k}{b} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & -\frac{3k}{b} & -\frac{3k}{b} & 0 & -\frac{(\lambda^2 \rho_2 - 3k)}{b} & 0 \end{pmatrix}}_{\mathcal{B}_\lambda} \underbrace{\begin{pmatrix} u \\ u_x \\ z \\ z_x \\ s \\ s_x \end{pmatrix}}_{\mathcal{W}}. \quad (4.12)$$

From (4.6), (4.7) and (4.9) and the regularity (Theorem 3.4), we have $\mathcal{W}(\alpha) = 0$. From (4.11) we have that the solution is given by

$$\mathcal{W}(x) = e^{\mathcal{B}_\lambda (x-\alpha)} \mathcal{W}(\alpha). \quad (4.13)$$

Thereby, from (4.13) and using that $\mathcal{W}(\alpha) = 0$, we get

$$\mathcal{W} = 0. \quad (4.14)$$

This way from (4.14) and using that $u(0) = z(0) = s(0) = 0$, it follows that

$$u = 0, \quad z = 0, \quad \text{and} \quad s = 0 \quad \text{in} \quad (0, \alpha). \quad (4.15)$$

Therefore from (4.15), (4.2)_{1,3,5} and the fact that $\lambda \neq 0$, we obtain

$$U = 0 \quad \text{in} \quad (0, \alpha). \quad (4.16)$$

From (4.15) and the regularity of z and s , we obtain

$$z(\alpha) = 0 \quad \text{and} \quad s(\alpha) = 0, \quad (4.17)$$

where using (4.9) and the same idea given above we obtain

$$u = 0, \quad z = 0 \quad \text{and} \quad s = 0 \quad \text{in} \quad (\alpha, \beta). \quad (4.18)$$

This way, from (4.2)_{1,3,5} and the fact that $\lambda \neq 0$, we obtain

$$U = 0 \quad \text{in} \quad (\alpha, \beta). \quad (4.19)$$

Now, let $\widetilde{\mathcal{W}} = (u, u_x, z, z_x, s, s_x)^T$. From (4.19) and the regularity of u , z and s we have $\widetilde{\mathcal{W}}(\beta) = 0$ and the system (4.10) in (β, L) implies

$$\widetilde{\mathcal{W}}_x = \mathcal{B}_\lambda \widetilde{\mathcal{W}} \quad \text{in} \quad (\beta, L), \quad (4.20)$$

where \mathcal{B}_λ is defined in (4.12). This way, we have

$$\widetilde{\mathcal{W}}(x) = e^{\mathcal{B}_\lambda (x-\beta)} \widetilde{\mathcal{W}}(\beta) = 0. \quad (4.21)$$

Therefore, (4.2)_{1,3,5} we conclude

$$U = 0 \quad \text{in} \quad (\beta, L). \quad (4.22)$$

Finally, from (4.16), (4.19) and (4.22) we obtain

$$U = 0 \quad \text{in} \quad (0, L). \quad (4.23)$$

The Lemma 4.1 follows. \square

Lemma 4.2. *For all $\lambda \in \mathbb{R}$, we have*

$$R(i\lambda I - \mathcal{A}) = \mathcal{H}. \quad (4.24)$$

Proof. By the Lemma 3.2, $0 \in \varrho(\mathcal{A})$. We will show the result for $\lambda \in \mathbb{R}^*$. For this aim, let $\mathcal{F} = (f_1, f_2, f_3, f_4, f_5, f_6)^T \in \mathcal{H}$, we want to find $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ solution of

$$(i\lambda I - \mathcal{A})W = \mathcal{F}, \quad (4.25)$$

that is

$$\begin{cases} i\lambda u - U = f_1, \\ i\lambda \rho_1 U + k(s - z - u_x)_x - k d(x)(S_x - U) = \rho_1 f_2, \\ i\lambda z - Z = f_3, \\ i\lambda \rho_2 Z - b z_{xx} - k(s - z - u_x) + \beta Z = \rho_2 f_4, \\ i\lambda s - S = f_5, \\ i\lambda \rho_2 S - b s_{xx} + 3k(s - z - u_x) - 3k[d(x)(S_x - U)]_x + 4\delta \rho_2 S = \rho_2 f_6, \end{cases} \quad (4.26)$$

with the boundary conditions

$$u(0) = u(L) = z(0) = z(L) = s(0) = s(L) = 0. \quad (4.27)$$

From (4.26)_{1,3,5} we have

$$\begin{cases} U = i\lambda u - f_1, \\ Z = i\lambda z - f_3, \\ S = i\lambda s - f_5. \end{cases} \quad (4.28)$$

Replacing (4.28)_{1,2,3} into (4.26)_{2,4,6} we obtain

$$-\lambda^2 u + \frac{k}{\rho_1} [(s - z - u_x)_x - i\lambda d(x)(s_x - u)] = g_1(x), \quad (4.29)$$

$$-\lambda^2 z - \frac{1}{\rho_2} [b z_{xx} + k(s - z - u_x)] + i\lambda \beta z = g_1(x), \quad (4.30)$$

$$-\lambda^2 s - \frac{1}{\rho_2} [b s_{xx} - 3k(s - z - u_x) + 3i\lambda k[d(x)(s_x - u)]_x] + 4i\lambda \delta s = g_3(x), \quad (4.31)$$

where

$$\begin{cases} g_1(x) = f_2 + i\lambda f_1 - \frac{k}{\rho_1} d(x)[f_{5x} - f_1] \in H^{-1}(0, L), \\ g_2(x) = f_4 + (i\lambda + b)f_3 \in H^{-1}(0, L), \\ g_3 = f_6 + i\lambda f_5 + 4\delta f_5 - \frac{3k}{\rho_2} [d(x)(f_{5x} - f_1)]_x \in H^{-1}(0, L). \end{cases} \quad (4.32)$$

On the other hand, for all $\widetilde{W} = (u, z, s)^T \in \widetilde{\mathcal{H}} = [H_0^1(0, L)]^3$, we define the linear operator

$$\widetilde{\mathcal{L}} : \widetilde{\mathcal{H}} \rightarrow \widetilde{\mathcal{H}}' := [H^{-1}(0, L)]^3 \quad (4.33)$$

by

$$\tilde{\mathcal{L}}(\tilde{W}) = \begin{pmatrix} \frac{k}{\rho_1} [(s-z-u_x)_x - i\lambda d(x)(s_x-u)] \\ -\frac{1}{\rho_2} [bz_{xx} + k(s-z-u_x)] + i\lambda bz \\ -\frac{1}{\rho_2} [bs_{xx} - 3k(s-z-u_x) + 3i\lambda k[d(x)(s_x-u)]_x] + 4i\lambda\delta s \end{pmatrix}. \quad (4.34)$$

Claim: $\tilde{\mathcal{L}}$ is an isomorphism. In fact, let $\Psi = (\rho_1 \Psi_1, \rho_2 \Psi_2, \rho_2 \Psi_3)^T \in \tilde{\mathcal{H}}$, then

$$\begin{aligned} \langle \tilde{\mathcal{L}}(\tilde{W}), \Psi \rangle_{\tilde{\mathcal{H}}' \times \tilde{\mathcal{H}}} &= \langle k(s-z-u_x)_x - i\lambda k d(x)(s_x-u), \Psi_1 \rangle_{H^{-1}(0,L) \times H_0^1(0,L)} \\ &\quad + \langle -bz_{xx} - k(s-z-u_x) + i\rho_2 \lambda bz, \Psi_2 \rangle_{H^{-1}(0,L) \times H_0^1(0,L)} \\ &\quad + \langle -bs_{xx} + 3k(s-z-u_x) - 3i\lambda k[d(x)(s_x-u)]_x \\ &\quad + 4i\rho_2 \lambda \delta s, \Psi_3 \rangle_{H^{-1}(0,L) \times H_0^1(0,L)}. \end{aligned}$$

Then is not difficult to show that

$$\begin{aligned} \langle \tilde{\mathcal{L}}(\tilde{W}), \Psi \rangle_{\tilde{\mathcal{H}}' \times \tilde{\mathcal{H}}} &= -k \int_0^L (s-z-u_x) \bar{\Psi}_{1x} dx - ik\lambda \int_0^L d(x)(s_x-u) \bar{\Psi}_1 dx \\ &\quad + b \int_0^L z_x \bar{\Psi}_{2x} dx - k \int_0^L (s-z-u_x) \bar{\Psi}_2 dx + i\rho_2 \lambda b \int_0^L z \bar{\Psi}_2 dx \\ &\quad + b \int_0^L s_x \bar{\Psi}_{3x} dx + 3k \int_0^L (s-z-u_x) \bar{\Psi}_3 dx \\ &\quad + 3i\lambda k \int_0^L d(x)(s_x-u) \bar{\Psi}_{3x} dx + 4i\lambda \rho_2 \delta \int_0^L s \bar{\Psi}_3 dx \end{aligned}$$

defines a continuous sesquilinear form which is coercive on \mathcal{H} . Therefore, using the Lax-Milgram theorem, we deduce that $\tilde{\mathcal{L}}$ is an isomorphism from $\tilde{\mathcal{H}}$ onto $\tilde{\mathcal{H}}'$. On the other hand, let $\tilde{W} = (u, z, s)^T$ and $\mathcal{G} = (g_1, g_2, g_3)^T$, then (4.29), (4.30) and (4.31) can be transformed into the following form

$$(I - \lambda^2 \tilde{\mathcal{L}}^{-1}) \tilde{W} = \tilde{\mathcal{L}}^{-1} \mathcal{G}. \quad (4.35)$$

Using $I : \tilde{\mathcal{H}} \rightarrow \tilde{\mathcal{H}}'$ is a compact operator and $\tilde{\mathcal{L}}^{-1} : \tilde{\mathcal{H}}' \rightarrow \tilde{\mathcal{H}}$ is an isomorphism, then the operator $I - \lambda^2 \tilde{\mathcal{L}}^{-1}$ is Fredholm of index zero. Then, by the Fredholm alternative, (4.35) admits a unique solution $\tilde{W} \in \tilde{\mathcal{H}}$ if and only if $I - \lambda^2 \tilde{\mathcal{L}}^{-1}$ is injective. Let $\mathbf{W} = (\mathbf{u}, \mathbf{z}, \mathbf{s})^T \in \tilde{\mathcal{H}}$ such that

$$\mathbf{W} - \lambda^2 \tilde{\mathcal{L}}^{-1} \mathbf{W} = 0 \iff \lambda^2 \mathbf{W} - \tilde{\mathcal{L}} \mathbf{W} = 0. \quad (4.36)$$

Equivalently, we have

$$-\lambda^2 \mathbf{u} + \frac{k}{\rho_1} [(s-z-u_x)_x - i\lambda d(x)(s_x-u)] = 0, \quad (4.37)$$

$$-\lambda^2 \mathbf{z} - \frac{1}{\rho_2} [bz_{xx} + k(s-z-u_x)] + i\lambda b \mathbf{z} = 0, \quad (4.38)$$

$$-\lambda^2 \mathbf{s} - \frac{1}{\rho_2} [bs_{xx} - 3k(s-z-u_x) + 3i\lambda k[d(x)(s_x-u)]_x] + 4i\lambda\delta \mathbf{s} = 0. \quad (4.39)$$

It is not difficult to show that if $\mathbf{W} = (\mathbf{u}, \mathbf{z}, \mathbf{s})^T$ is a solution of (4.37)-(4.39), then $\widetilde{\mathbf{W}} = (\mathbf{u}, i\lambda\mathbf{u}, \mathbf{z}, i\lambda\mathbf{z}, \mathbf{s}, i\lambda\mathbf{s})^T \in \mathcal{D}(\mathcal{A})$ and satisfies: $i\lambda\mathbf{W} - \mathcal{A}\mathbf{W} = 0$. Thereby, using Lemma 4.1 we obtain $\widetilde{\mathbf{W}} = 0$ and we conclude that $I - \lambda^2 \widetilde{\mathcal{L}}^{-1}$ is injective. Since to Fredholm's alternative, (4.35) admits a unique solution $\widetilde{W} \in \mathcal{H}$ and

$$u, z, s \in H^2(0, L), \quad 3kd(x)[i\lambda s_x - f_{5x} - (i\lambda u - f_1)]_x \in L^2(0, L).$$

Finally, setting $U = i\lambda u - f_1$, $Z = i\lambda z - f_3$, and $S = i\lambda s - f_5$ we have that $W \in \mathcal{D}(\mathcal{A})$ is a unique solution of (4.25). The Lemma 4.2 follows. \square

Now we will prove the main theorem of this section:

Theorem 4.3. *The C_0 -semigroup of contraction $(e^{t\mathcal{A}})_{t \geq 0}$ is strongly stable in \mathcal{H} , that is, for all $W_0 \in \mathcal{H}$, the solution of (3.3) satisfies*

$$\lim_{t \rightarrow +\infty} \|e^{t\mathcal{A}}W_0\|_{\mathcal{H}} = 0.$$

Proof. We have that the operator \mathcal{A} has no pure imaginary, that is, $\sigma_p(\mathcal{A}) \cap i\mathbb{R} = \emptyset$ (Lemma 4.1). On the other hand, using the closed graph theorem of Banach and Lemma 4.2, we have $\sigma(\mathcal{A}) \cap i\mathbb{R} = \emptyset$. Now using Theorem 2.3 due to Arendt-Batty [3], we get that the C_0 -semigroup $(e^{t\mathcal{A}})_{t \geq 0}$ is strongly stable. The main theorem follows. \square

5. POLYNOMIAL STABILITY

In this section, we will prove the polynomial stability of the system (1.1). Since $i\mathbb{R} \subset \varrho(\mathcal{A})$ (Section three), from Theorem 2.4 to prove the main theorem, we still need to prove the following condition

$$\sup_{\lambda \in \mathbb{R}} \left\| (i\lambda I - \mathcal{A})^{-1} \right\|_{\mathcal{L}(\mathcal{H})} = \mathcal{O}(|\lambda|^\ell). \quad (5.1)$$

We will prove condition (5.1) by a contradiction argument. For this purpose, we suppose that (5.1) is false, then there exists

$$\left[\left(\lambda^{(n)}, W^{(n)} = \left(u^{(n)}, U^{(n)}, z^{(n)}, Z^{(n)}, s^{(n)}, S^{(n)} \right)^T \right) \right]_{n \geq 1} \subseteq \mathbb{R}^* \times \mathcal{D}(\mathcal{A})$$

with

$$|\lambda^{(n)}| \rightarrow \infty \quad \text{and} \quad \left\| W^{(n)} \right\|_{\mathcal{H}} = \left\| \left(u^{(n)}, U^{(n)}, z^{(n)}, Z^{(n)}, s^{(n)}, S^{(n)} \right)^T \right\|_{\mathcal{H}} = 1, \quad (5.2)$$

such that

$$\left(\lambda^{(n)} \right)^\ell \left(i\lambda^{(n)} I - \mathcal{A} \right) W^{(n)} = F^{(n)} = (f_1^{(n)}, f_2^{(n)}, f_3^{(n)}, f_4^{(n)}, f_5^{(n)}, f_6^{(n)}) \rightarrow 0. \quad (5.3)$$

Here for simplicity, we omit the index n . Thereby from (5.3) we have

$$\begin{cases} i \lambda u - U = \lambda^{-\ell} f_1 \iff U = i \lambda u - \lambda^{-\ell} f_1, \\ i \lambda \rho_1 U + k(s - z - u_x)_x - k d(x)(S_x - U) = \rho_1 \lambda^{-\ell} f_2, \\ i \lambda z - Z = \lambda^{-\ell} f_3 \iff Z = i \lambda z - \lambda^{-\ell} f_3, \\ i \lambda \rho_2 Z - b z_{xx} - k(s - z - u_x) + \beta Z = \rho_2 \lambda^{-\ell} f_4, \\ i \lambda s - S = \lambda^{-\ell} f_5 \iff S = i \lambda s - \lambda^{-\ell} f_5, \\ i \lambda \rho_2 S - b s_{xx} + 3k(s - z - u_x) - 3k[d(x)(S_x - U)]_x + 4\delta \rho_2 S = \rho_2 \lambda^{-\ell} f_6. \end{cases} \quad (5.4)$$

Replacing (5.4)_{1, 3, 5} into (5.4)_{2, 4, 6} respectively and performing straightforward calculations we obtain

$$\lambda^2 \rho_1 u - k(s - z - u_x)_x + k d(x)(S_x - U) = -\rho_1 \lambda^{-\ell} f_2 - i \rho_1 \lambda^{-\ell+1} f_1, \quad (5.5)$$

$$\lambda^2 \rho_2 z + b z_{xx} + k(s - z - u_x) - \beta Z = -\rho_2 \lambda^{-\ell} f_4 - i \rho_2 \lambda^{-\ell+1} f_3, \quad (5.6)$$

$$\begin{aligned} \lambda^2 \rho_2 s + b s_{xx} - 3k(s - z - u_x) + 3k[d(x)(S_x - U)]_x + 4\delta S \\ = -\rho_2 \lambda^{-\ell} f_6 - i \rho_2 \lambda^{-\ell+1} f_5. \end{aligned} \quad (5.7)$$

Here we will check the condition (5.1) by finding a contradiction with (5.2) by proving that $\|W\|_{\mathcal{H}} = \mathcal{O}(1)$. From (5.4), and the fact that $\ell = 2$,

$$\|W\|_{\mathcal{H}} = 1 \quad \text{and} \quad \|F\|_{\mathcal{H}} \approx \mathcal{O}(1), \quad (5.8)$$

we have that

$$\begin{cases} \|u\| \approx \mathcal{O}(|\lambda|^{-1}), \quad \|z\| \approx \mathcal{O}(|\lambda|^{-1}) \quad \text{and} \quad \|s\| \approx \mathcal{O}(|\lambda|^{-1}), \\ \|u_{xx}\| \approx \mathcal{O}(|\lambda|), \quad \|z_{xx}\| \approx \mathcal{O}(|\lambda|) \quad \text{and} \quad \|s_{xx}\| \approx \mathcal{O}(|\lambda|), \\ [d(x)(S_x - U)]_x \approx \mathcal{O}(|\lambda|). \end{cases} \quad (5.9)$$

Moreover, from Poincaré inequality and the fact $\|F\| = \mathcal{O}(1)$ it follows that

$$\|f_1\| \lesssim \|f_{1x}\| \approx \mathcal{O}(1), \quad \|f_3\| \lesssim \|f_{3x}\| \approx \mathcal{O}(1) \quad \text{and} \quad \|f_5\| \lesssim \|f_{5x}\| \approx \mathcal{O}(1). \quad (5.10)$$

Claim 5.1. *Let $\ell = 2$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies*

$$\begin{cases} \int_{\alpha}^{\beta} |S_x - U|^2 dx \lesssim \frac{\mathcal{O}(1)}{\lambda^{\ell}}, \quad \|S_x - U\|^2 \lesssim \frac{\mathcal{O}(1)}{\lambda^{\ell}}, \quad \|S_x - U\| \lesssim \frac{\mathcal{O}(1)}{\lambda^{\ell/2}}, \\ \int_{\alpha}^{\beta} |s_x - u|^2 dx \lesssim \frac{\mathcal{O}(1)}{\lambda^{\ell+2}}, \\ \int_{\alpha}^{\beta} |S_x|^2 dx \lesssim \mathcal{O}(1), \\ \int_{\alpha}^{\beta} |s_x|^2 dx \lesssim \frac{\mathcal{O}(1)}{\lambda^2}. \end{cases} \quad (5.11)$$

Proof. Taking the inner product of (5.3) with W we have

$$- \operatorname{Re} \langle \mathcal{A}W, W \rangle_{\mathcal{H}} = \lambda^{-\ell} \operatorname{Re} \langle F, W \rangle_{\mathcal{H}} \leq \lambda^{-\ell} \|F\|_{\mathcal{H}} \|W\|_{\mathcal{H}}.$$

Then taking the real part and using (3.6) we obtain

$$\beta \int_0^L |Z|^2 dx + 4\delta \int_0^L |S|^2 dx + 3k d_0 \int_\alpha^\beta |S_x - U|^2 dx \leq \lambda^{-\ell} \underbrace{\|F\|_{\mathcal{H}}}_{\lesssim \mathcal{O}(1)} \underbrace{\|W\|_{\mathcal{H}}}_{=1}. \quad (5.12)$$

From (5.8) we obtain

$$\beta \int_0^L |Z|^2 dx + 4\delta \int_0^L |S|^2 dx + 3k d_0 \int_\alpha^\beta |S_x - U|^2 dx \lesssim \frac{\mathcal{O}(1)}{\lambda^\ell}. \quad (5.13)$$

From (5.13) we have that (5.11)₁ follows. Differentiating (5.4)₅ in x -variable we have

$$i \lambda s_x - S_x = \lambda^{-\ell} f_{5x}. \quad (5.14)$$

Subtracting (5.4)₁ with (5.14) we obtain

$$i \lambda (s_x - u) = (S_x - U) + \lambda^{-\ell} (f_{5x} - f_1).$$

Then using $(a + b)^2 \leq 2a^2 + 2b^2$ we have

$$\begin{aligned} \int_\alpha^\beta |s_x - u|^2 dx &\leq \frac{2}{\lambda^2} \int_\alpha^\beta |S_x - U|^2 dx + \frac{2}{\lambda^{2\ell+2}} \int_\alpha^\beta |f_{5x} - f_1|^2 dx \\ &\leq \frac{2}{\lambda^2} \int_\alpha^\beta |S_x - U|^2 dx + \frac{4}{\lambda^{2\ell+2}} \int_\alpha^\beta |f_{5x}|^2 dx + \frac{4}{\lambda^{2\ell+2}} \int_\alpha^\beta |f_1|^2 dx \\ &\leq \frac{2}{\lambda^2} \int_\alpha^\beta |S_x - U|^2 dx + \frac{4}{\lambda^{2\ell+2}} \|f_{5x}\|^2 + \frac{4}{\lambda^{2\ell+2}} \|f_1\|^2. \end{aligned}$$

Now using (5.10) together with we obtain

$$\int_\alpha^\beta |s_x - u|^2 dx \approx \frac{2}{\lambda^2} \frac{\mathcal{O}(1)}{\lambda^{2\ell}} + \frac{4}{\lambda^{2\ell+2}} \mathcal{O}(1) + \frac{4}{\lambda^{2\ell+2}} \mathcal{O}(1) \lesssim \frac{1}{\lambda^{2\ell+2}} \mathcal{O}(1), \quad (5.15)$$

then (5.11)₂ follows. On the other hand, using (5.13)

$$\begin{aligned} \int_\alpha^\beta |S_x|^2 dx &= \int_\alpha^\beta |S_x - U + U|^2 dx \\ &\leq 2 \int_\alpha^\beta |S_x - U|^2 dx + 2 \int_\alpha^\beta |U|^2 dx \\ &\lesssim \frac{2}{\lambda^\ell} \mathcal{O}(1) + 2 \int_\alpha^\beta |U|^2 dx. \end{aligned} \quad (5.16)$$

Using that U is uniformly bounded in $L^2(0, L)$, we get (5.11)₃. Differentiating (5.4)₅ we have

$$i \lambda s_x = S_x + \lambda^{-\ell} f_{5x}.$$

Elevating square the above expression, integrating over (α, β) and using $(a + b)^2 \leq 2a^2 + 2b^2$ together with (5.11)₃ and (5.10) we obtain

$$\begin{aligned} \int_{\alpha}^{\beta} |s_x|^2 dx &\leq \frac{2}{\lambda^2} \int_{\alpha}^{\beta} |S_x|^2 dx + \frac{2}{\lambda^{2\ell+2}} \int_{\alpha}^{\beta} |f_{5x}|^2 dx \\ &\lesssim \frac{2}{\lambda^2} \mathcal{O}(1) + \frac{2}{\lambda^{2\ell+2}} \mathcal{O}(1) \\ &\lesssim \left(\frac{2}{\lambda^2} + \frac{2}{\lambda^2} \right) \mathcal{O}(1) \lesssim \frac{4}{\lambda^2} \mathcal{O}(1) \approx \frac{C}{\lambda^2} \mathcal{O}(1), \end{aligned} \quad (5.17)$$

then we have (5.11)₄. The Claim 5.1 follows. \square

For all $0 < \varepsilon < \frac{(\beta-\alpha)}{10}$, we fix the following cut-off functions:

$\chi_j \in C^\infty([0, L])$, $j \in \{1, \dots, 4\}$ such that $0 \leq \chi_j \leq 1$, for all $x \in [0, L]$ and

$$\chi_j(x) = \begin{cases} 1 & \text{if } x \in [\alpha + j\varepsilon, \beta - j\varepsilon], \\ 0 & \text{if } x \in [0, \alpha + (j-1)\varepsilon] \cup [\beta + (1-j)\varepsilon, L]. \end{cases} \quad (5.18)$$

Claim 5.2. *Let $\ell = 2$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies the following estimates*

$$\begin{cases} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |S|^2 dx \approx \mathcal{O}(1), \\ \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |\lambda s|^2 dx \approx \mathcal{O}(1). \end{cases} \quad (5.19)$$

Proof. Multiplying (5.4)₆ by $-i\lambda^{-1}\chi_1(x)\bar{S}$ and integrating over (α, β) we have

$$\begin{aligned} \rho_2 \int_{\alpha}^{\beta} \chi_1(x) |S|^2 dx &= -\frac{ib}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) s_{xx} \bar{S} dx + \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) (s - z - u_x) \bar{S} dx \\ &\quad - \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) [d(x)(S_x - U)]_x \bar{S} dx \\ &\quad + \frac{4i\delta\rho_2}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) |S|^2 dx - \frac{i\rho_2}{\lambda^{\ell+1}} \int_{\alpha}^{\beta} \chi_1(x) f_6 \bar{S} dx. \end{aligned} \quad (5.20)$$

Using the Cauchy-Schwarz inequality and the fact that S is uniformly bounded in $L^2(0, L)$ together with (5.10) and $\chi_1(\alpha) = \chi_1(\beta) = 0$ we have

$$\begin{aligned} \rho_2 \int_{\alpha}^{\beta} \chi_1(x) |S|^2 dx &\approx \frac{ib}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) s_x \bar{S}_x dx + \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) (s - z - u_x) \bar{S} dx \\ &\quad - \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) [d(x)(S_x - U)]_x \bar{S} dx + \frac{\mathcal{O}(1)}{\lambda^{\ell+1}}. \end{aligned} \quad (5.21)$$

Moreover, using the fact that $(s - z - u_x)$ and S are uniformly bounded in $L^2(0, L)$, together with Claim 5.1 we get

$$\rho_2 \int_{\alpha}^{\beta} \chi_1(x) |S|^2 dx \approx \underbrace{-\frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) [d(x)(S_x - U)]_x \bar{S} dx}_{I_1} + \mathcal{O}(1). \quad (5.22)$$

Integrating by parts, using the definition of $d(x)$ and the fact that $\chi_1(\alpha) = \chi_1(\beta) = 0$ it follows that

$$\begin{aligned} I_1 &= \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_{1x}(x) d(x) (S_x - U) \bar{S} dx + \frac{3ik}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) d(x) (S_x - U) \bar{S}_x dx \\ &= \frac{3ikd_0}{\lambda} \int_{\alpha}^{\beta} \chi_{1x}(x) (S_x - U) \bar{S} dx + \frac{3ikd_0}{\lambda} \int_{\alpha}^{\beta} \chi_1(x) (S_x - U) \bar{S}_x dx. \end{aligned} \quad (5.23)$$

From Claim 5.1 and the fact that S is uniformly bounded in $L^2(0, L)$ we obtain

$$I_1 \approx \frac{\mathcal{O}(1)}{\lambda^{\frac{\ell}{2}+1}}. \quad (5.24)$$

Replacing (5.24) into (5.22) we obtain

$$\rho_2 \int_{\alpha}^{\beta} \chi_1(x) |S|^2 dx \approx \mathcal{O}(1). \quad (5.25)$$

From (5.25) and the definition (5.19), we have (5.18)₁. On the other hand, from (5.4)₅ and using that $(P + Q)^2 \leq 2P^2 + 2Q^2$ we have

$$|\lambda s|^2 = |i\lambda s|^2 = |S + \lambda^{-\ell} f_5|^2 \leq 2|S|^2 + 2|\lambda^{-\ell} f_5|^2 = 2|S|^2 + 2\lambda^{-2\ell} |f_5|^2.$$

Then

$$\begin{aligned} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |\lambda s|^2 dx &\leq 2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |S|^2 dx + 2\lambda^{-2\ell} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |f_5|^2 dx \\ &\approx 2\mathcal{O}(1) + 2\lambda^{-2\ell} \mathcal{O}(1) \\ &\lesssim 2\mathcal{O}(1). \end{aligned} \quad (5.26)$$

From (5.19)₁ we have the first estimate in the right-hand side and for the second term we use the fact that $\|f_5\| \approx \mathcal{O}(1)$. Then we obtain (5.19)₂. The Claim 5.2 follows. \square

Claim 5.3. *Let $\ell = 2$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies the following estimates*

$$\begin{cases} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |u_x|^2 dx \approx \mathcal{O}(1), \\ \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |\lambda u|^2 dx \approx \mathcal{O}(1). \end{cases} \quad (5.27)$$

Proof. Multiplying (5.4)₆ by $\chi_2(x) \bar{u}_x$ integrating over $(\alpha + \varepsilon, \beta - \varepsilon)$ and using the fact that u_x is uniformly bounded in $L^2(0, L)$ and $\|f_6\| \approx \mathcal{O}(1)$ we have

$$\begin{aligned}
 & \underbrace{i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S \bar{u}_x dx}_{J_1} - \underbrace{3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) [d(x) (S_x - U)]_x \bar{u}_x dx}_{J_2} \\
 & - 3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |u_x|^2 dx + 3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (s - z) \bar{u}_x dx \\
 & - b \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) s_{xx} \bar{u}_x dx + 4\delta \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S \bar{u}_x dx \\
 & = \rho_2 \lambda^{-\ell} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) f_6 \bar{u}_x dx \lesssim C \lambda^{-\ell} \mathcal{O}(1) \approx \frac{\mathcal{O}(1)}{\lambda^\ell}. \tag{5.28}
 \end{aligned}$$

Using the fact that u_x is uniformly bounded in $L^2(0, L)$, and from (5.9), $\|z\| \approx \mathcal{O}(|\lambda|^{-1})$ and $\|s\| \approx \mathcal{O}(|\lambda|^{-1})$ we get

$$\begin{cases} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (s - z) \bar{u}_x dx \approx \mathcal{O}(1), \\ \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) s_{xx} \bar{u}_x dx \approx \mathcal{O}(1), \\ \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S \bar{u}_x dx \approx \mathcal{O}(1). \end{cases} \tag{5.29}$$

Then

$$\begin{aligned}
 & \underbrace{-i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S \bar{u}_x dx}_{J_1} + \underbrace{3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) [d(x) (S_x - U)]_x \bar{u}_x dx}_{J_2} \\
 & + 3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |u_x|^2 dx \approx \mathcal{O}(1). \tag{5.30}
 \end{aligned}$$

On the other hand, using integration by parts, the definition of χ_2 , the Claim 5.2 ((5.19)₁), and the fact that $\|u\| \approx \mathcal{O}(|\lambda|^{-\ell})$ we obtain

$$\begin{aligned}
 J_1 &= -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx - i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_{2x}(x) S \bar{u} dx \\
 &\approx -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx + C \mathcal{O}(|\lambda|^{-\ell}) \mathcal{O}(1) \\
 &\lesssim -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx + \mathcal{O}(1). \tag{5.31}
 \end{aligned}$$

Moreover

$$\begin{aligned}
 & -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx = -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (S_x - U + U) \bar{u} dx \\
 & = -i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (S_x - U) \bar{u} dx - i \lambda \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) U \bar{u} dx. \tag{5.32}
 \end{aligned}$$

From Claim 5.1 (see (5.11)₁) and the fact that $\|u\| \approx \mathcal{O}(|\lambda|^{-1})$ we have

$$-i\lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx = -i\lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) U \bar{u} dx + \mathcal{O}(|\lambda|^{-\ell/2}). \quad (5.33)$$

Replacing (5.4)₁, that is, $U = i\lambda u - \lambda^{-\ell} f_1$ into (5.33) we have

$$\begin{aligned} -i\lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx &= -i\lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (i\lambda u - \lambda^{-\ell} f_1) \bar{u} dx + \mathcal{O}(|\lambda|^{-\ell/2}) \\ &= \lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |\lambda u|^2 dx + i\lambda^{-\ell+1} \rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) f_1 \bar{u} dx + \mathcal{O}(|\lambda|^{-\ell/2}). \end{aligned}$$

Now, using $\|u\| \approx \mathcal{O}(|\lambda|^{-1})$ and $\|f_1\| \approx \mathcal{O}(1)$ we obtain

$$-i\lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) S_x \bar{u} dx \approx \lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |\lambda u|^2 dx + \mathcal{O}(|\lambda|^{-\ell/2}). \quad (5.34)$$

Replacing (5.34) into (5.31) we obtain

$$J_1 \approx \lambda\rho_2 \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |\lambda u|^2 dx + \mathcal{O}(1). \quad (5.35)$$

On the other hand, using integration by parts and the definition of $\chi_2(x)$ we have

$$\begin{aligned} J_2 &= -3k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) [d(x) (S_x - U)]_x \bar{u}_x dx \\ &= 3d_0 k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_{2x}(x) d_0 (S_x - U) \bar{u}_x dx + 3d_0 k \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) (S_x - U) \bar{u}_{xx} dx. \end{aligned}$$

From Claim 5.1 and the fact that u_x is uniformly bounded in $L^2(0, L)$, $\|u_{xx}\| \approx \mathcal{O}(|\lambda|)$ (see (5.9)) we obtain

$$J_2 \approx \mathcal{O}(|\lambda|^{-\ell/2}). \quad (5.36)$$

Replacing (5.35) and (5.36) into (5.30)

$$\int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |u_x|^2 dx \approx \mathcal{O}(1) \quad \text{and} \quad \int_{\alpha+\varepsilon}^{\beta-\varepsilon} \chi_2(x) |\lambda u|^2 dx \approx \mathcal{O}(1). \quad (5.37)$$

Finally, from the above estimation and the definition of $\chi_2(x)$, we obtain (5.30). The Claim 5.3 follows. \square

Claim 5.4. Let $\ell = 2$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies the following estimates

$$\begin{cases} \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |z_x|^2 dx \approx \mathcal{O}(1), \\ \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |\lambda z|^2 dx \approx \mathcal{O}(1). \end{cases} \quad (5.38)$$

Proof. Multiplying (5.5) by $\rho_1^{-1} \chi_3(x) \bar{z}_x$, integrating over $(\alpha + 2\varepsilon, \beta - 2\varepsilon)$ and using the definition of $d(x)$ together with the fact that $z z_x$ is uniformly bounded

in $L^2(0, L)$, $\|f_1\| \approx \mathcal{O}(1)$ and $\|f_2\| \approx \mathcal{O}(1)$ we have

$$\begin{aligned} \lambda^2 \rho_1 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx - \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (s - z - u_x)_x \bar{z}_x dx \\ + \frac{k d_0}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (S_x - U) \bar{z}_x dx \approx \mathcal{O}(|\lambda|^{-1}) \end{aligned}$$

or

$$\begin{aligned} \lambda^2 \rho_1 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx - \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) s_x \bar{z}_x dx \\ + \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) |z_x|^2 dx + \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_{xx} \bar{z}_x dx \\ + \frac{k d_0}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (S_x - U) \bar{z}_x dx \approx \mathcal{O}(|\lambda|^{-1}). \end{aligned}$$

Then

$$\begin{aligned} \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) |z_x|^2 dx \lesssim -\lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx + \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) s_x \bar{z}_x dx \\ - \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_{xx} \bar{z}_x dx - \frac{k d_0}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (S_x - U) \bar{z}_x dx + \mathcal{O}(|\lambda|^{-1}). \end{aligned} \quad (5.39)$$

From Claim 5.1, (5.18) ($j = 3$) and the fact that z_x is uniformly bounded in $L^2(0, L)$, it follows that ($\ell = 2$)

$$\begin{cases} \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) s_x \bar{z}_x dx \approx \mathcal{O}(1) \\ -\frac{k d_0}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (S_x - U) \bar{z}_x dx \approx \mathcal{O}(|\lambda|^{-\ell+1}) \approx \mathcal{O}(|\lambda|^{-1}). \end{cases} \quad (5.40)$$

Replacing (5.40) into (5.39)

$$\frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) |z_x|^2 dx \lesssim -\lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx \quad (5.41)$$

$$- \frac{k}{\rho_1} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_{xx} \bar{z}_x dx + \mathcal{O}(1). \quad (5.42)$$

On the other hand, from (5.6) and applying conjugate we have ($\ell = 2$)

$$\lambda^2 \rho_2 \bar{z} + b \bar{z}_{xx} + k (\bar{s} - \bar{z} - \bar{u}_x) - \beta \bar{Z} = -\rho_2 \lambda^{-2} \bar{f}_4 + i \rho_2 \lambda^{-1} \bar{f}_3. \quad (5.43)$$

Multiplying (5.43) by $\rho_2^{-1} \chi_3 u_x$, integrating over $(\alpha + 2\varepsilon, \beta - 2\varepsilon)$ we obtain

$$\begin{aligned} \lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) \bar{z} u_x dx + \frac{b}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) \bar{z}_{xx} u_x dx \\ + \frac{k}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) (\bar{s} - \bar{z} - \bar{u}_x) u_x dx - \frac{\beta}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) \bar{Z} u_x dx \approx \mathcal{O}(|\lambda|^{-1}). \end{aligned} \quad (5.44)$$

Integrating the first two terms and the fourth term in (5.44) we obtain

$$\begin{aligned}
& - \lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx - \frac{b}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_{xx} \bar{z}_x dx \\
& + \frac{\beta}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{Z}_x dx \\
& \lesssim \lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_{3x}(x) u \bar{z} dx + \frac{b}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_{3x}(x) u_x \bar{z}_x dx \\
& + \frac{\beta}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_{3x}(x) u \bar{Z} dx + \frac{k}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_x (\bar{s} - \bar{z} - \bar{u}_x) dx + \mathcal{O}(|\lambda|^{-1}).
\end{aligned} \tag{5.45}$$

From Claim 5.3, using the fact that z_x , Z and $(\bar{s} - \bar{z} - \bar{u}_x)$ are uniformly bounded in $L^2(0, L)$ together with $\|z\| \approx \mathcal{O}(|\lambda|^{-\ell+1})$, we get

$$\begin{cases} \lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_{3x}(x) u \bar{z} dx \approx \mathcal{O}(1), \\ \frac{b}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_{3x}(x) u_x \bar{z}_x dx \approx \mathcal{O}(1), \\ \frac{k}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_x (\bar{s} - \bar{z} - \bar{u}_x) dx \approx \mathcal{O}(1). \end{cases} \tag{5.46}$$

Inserting (5.46) into (5.45) we get

$$- \lambda^2 \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u \bar{z}_x dx - \frac{b}{\rho_2} \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} \chi_3(x) u_{xx} \bar{z}_x dx \approx \mathcal{O}(1). \tag{5.47}$$

Replacing (5.47) into (5.42) and using the definition of χ_3 , we obtain (5.38)₁. On the other hand, multiplying (5.43) by $\chi_4(x) z$, integrating by parts over $(\alpha+3\varepsilon, \beta-3\varepsilon)$, using the definition of χ_4 and the fact $\|z\| = \mathcal{O}(1)$ together with $\|f_3\| = \mathcal{O}(1)$ and $\|f_4\| = \mathcal{O}(1)$ we get

$$\begin{aligned}
\rho_2 \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) |\lambda z|^2 dx & \lesssim b \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) |\bar{z}_x|^2 dx + \beta \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) Z z dx \\
& - k \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) (\bar{s} - \bar{z} - \bar{u}_x) z dx + \mathcal{O}(|\lambda|^{-2})
\end{aligned}$$

or

$$\begin{aligned}
\rho_2 \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) |\lambda z|^2 dx & \lesssim b \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) |\bar{z}_x|^2 dx \\
& - k \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) z (\bar{s} - \bar{z} - \bar{u}_x) dx + \mathcal{O}(|\lambda|^{-2}).
\end{aligned} \tag{5.48}$$

From (5.48), the first estimation in (5.38) and using the fact that $(\bar{s} - \bar{z} - \bar{u}_x)$ is uniformly bounded in $L^2(0, L)$ together with $\|z\| \approx \mathcal{O}(|\lambda|^{-1})$ we obtain

$$\int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} \chi_4(x) |\lambda z|^2 dx \approx \mathcal{O}(1). \tag{5.49}$$

Thereby, from (5.49) and the definition of $\chi_4(x)$, we obtain the second estimative. \square

Claim 5.5. Let $\ell = 2$. Let $F \in C^1([0, L])$ such that $F(0) = F(L) = 0$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies the following estimates

$$\int_0^L F_x [\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2 + 3k |d(x)(S_x - U)|^2] dx \approx \mathcal{O}(1). \quad (5.50)$$

Proof. First, we multiply (5.5) by $2F \bar{u}_x$ integrating over $(0, L)$. Later we use the Claim 5.1 and then we use the fact that u_x is uniformly bounded in $L^2(0, L)$ together with $\|u\| \approx \mathcal{O}(|\lambda|^{-1})$, $\|f_1\| \approx \mathcal{O}(1)$, $\|f_1\| \approx \mathcal{O}(1)$.

$$\lambda^2 \rho_1 \int_0^L 2F u \bar{u}_x dx - k \int_0^L 2F (s - z - u_x)_x \bar{u}_x dx + k \int_0^L 2F d(x)(S_x - U) \bar{u}_x dx \approx \mathcal{O}(|\lambda|^{-\ell+1}).$$

Now we take the real part, that is,

$$\begin{aligned} & \operatorname{Re} \left\{ \lambda^2 \rho_1 \int_0^L 2F u \bar{u}_x dx \right\} + \operatorname{Re} \left\{ k \int_0^L 2F u_{xx} \bar{u}_x dx \right\} \\ & + \operatorname{Re} \left\{ k \int_0^L 2F z_x \bar{u}_x dx \right\} - \operatorname{Re} \left\{ k \int_0^L 2F s_x \bar{u}_x dx \right\} \\ & + \operatorname{Re} \left\{ k \int_0^L 2F d(x)(S_x - U) \bar{u}_x dx \right\} \lesssim \mathcal{O}(|\lambda|^{-\ell+1}). \end{aligned} \quad (5.51)$$

Using that $2 \operatorname{Re}(h \bar{h}_x) = (|h|^2)_x$ we have

$$\begin{aligned} & \int_0^L F (\rho_1 |\lambda u|^2 + k |u_x|^2)_x dx + \operatorname{Re} \left\{ 2k \int_0^L F z_x \bar{u}_x dx \right\} \\ & - \operatorname{Re} \left\{ 2k \int_0^L F s_x \bar{u}_x dx \right\} + \underbrace{\operatorname{Re} \left\{ 2k d_0 \int_0^L F (S_x - U) \bar{u}_x dx \right\}}_{\approx \mathcal{O}(|\lambda|^{-\ell/2})} \lesssim \mathcal{O}(|\lambda|^{-\ell+1}). \end{aligned} \quad (5.52)$$

Multiplying (5.6) by $2F \bar{z}_x$ and we integrate over $(0, L)$. Later we use the fact that z_x is uniformly bounded in $L^2(0, L)$ together with $\|z\| \approx \mathcal{O}(|\lambda|^{-1})$, $\|s\| \approx \mathcal{O}(|\lambda|^{-1})$, $\|f_3\| \approx \mathcal{O}(1)$ and $\|f_4\| \approx \mathcal{O}(1)$, $2 \operatorname{Re}(h \bar{h}_x) = (|h|^2)_x$ we have

$$\begin{aligned} & \int_0^L F (\rho_2 |\lambda z|^2 + b |z_x|^2)_x dx - \operatorname{Re} \left\{ 2k \int_0^L F u_x \bar{z}_x dx \right\} \\ & + \underbrace{\operatorname{Re} \left\{ 2k \int_0^L F (s - z) \bar{z}_x dx \right\}}_{\approx \mathcal{O}(1)} \lesssim \mathcal{O}(|\lambda|^{1-\ell}). \end{aligned} \quad (5.53)$$

On the other hand, let $\mathcal{M} = d(x)(S_x - U)$. From Claim 5.1, the definition of $d(x)$ and the fact that s_x is uniformly bounded in $L^2(0, L)$, we get \mathcal{M} is uniformly bounded in $L^2(0, L)$. Multiplying (5.7) by $2F \bar{\mathcal{M}}$ and we integrate over $(0, L)$. Taking the

real part and using the fact that $\|z\| \approx \mathcal{O}(|\lambda|^{-1})$, $\|S\| \approx \mathcal{O}(|\lambda|^{-1})$, $\|f_5\| \approx \mathcal{O}(1)$ and $\|f_6\| \approx \mathcal{O}(1)$ we obtain

$$\begin{aligned} & \operatorname{Re} \left\{ 2 \lambda^2 \rho_2 \int_0^L F s \overline{\mathcal{M}} dx \right\} + \operatorname{Re} \left\{ 2b \int_0^L F s_{xx} \overline{\mathcal{M}} dx \right\} \\ & + \operatorname{Re} \left\{ -6k \int_0^L F (s - z - u_x) \overline{\mathcal{M}} dx \right\} \\ & + \operatorname{Re} \left\{ -6k \int_0^L F \mathcal{M}_x \overline{\mathcal{M}} dx \right\} + \operatorname{Re} \left\{ 8\delta \int_0^L F S \overline{\mathcal{M}} dx \right\} \\ & = \underbrace{\operatorname{Re} \left\{ -2 \int_0^L F (\rho_2 \lambda^{-\ell} f_6 - i \rho_2 \lambda^{-\ell+1} f_5) dx \right\}}_{\approx \mathcal{O}(|\lambda|^{-\ell+1})}. \end{aligned}$$

Using that $2 \operatorname{Re}(h \overline{h_x}) = (|h|^2)_x$ we have

$$\begin{aligned} & \operatorname{Re} \left\{ 2 \lambda^2 \rho_2 \int_0^L F s \overline{\mathcal{M}} dx \right\} + \operatorname{Re} \left\{ 2b \int_0^L F s_{xx} \overline{\mathcal{M}} dx \right\} \\ & + \operatorname{Re} \left\{ 6k \int_0^L F u_x \overline{\mathcal{M}} dx \right\} - \underbrace{\operatorname{Re} \left\{ 6k \int_0^L F (s - z) \overline{\mathcal{M}} dx \right\}}_{\approx \mathcal{O}(1)} \\ & - 3k \int_0^L F (|\mathcal{M}|^2)_x dx + \operatorname{Re} \left\{ 8\delta \int_0^L F S \overline{\mathcal{M}} dx \right\} \\ & = \underbrace{\operatorname{Re} \left\{ -2 \int_0^L F (\rho_2 \lambda^{-\ell} f_6 - i \rho_2 \lambda^{-\ell+1} f_5) \overline{\mathcal{M}} dx \right\}}_{\approx \mathcal{O}(|\lambda|^{-\ell+1})}. \quad (5.54) \end{aligned}$$

On the other hand, using the definition of $\mathcal{M} = d(x)(S_x - U)$ and $d(x)$, Claim 5.1 and the fact that u_x is uniformly bounded in $L^2(0, L)$, $\|s\| \approx \mathcal{O}(|\lambda|^{-1})$ we have

$$\operatorname{Re} \left\{ 2 \lambda^2 \rho_2 \int_0^L F s \overline{\mathcal{M}} dx \right\} = \underbrace{\operatorname{Re} \left\{ 2 \lambda^2 \rho_2 d_0 \int_0^L F s (\overline{S}_x - \overline{U}) dx \right\}}_{\approx \mathcal{O}(|\lambda|^{-\ell/2+1})} \quad (5.55)$$

and

$$\operatorname{Re} \left\{ 6k \int_0^L F u_x \overline{\mathcal{M}} dx \right\} = \underbrace{\operatorname{Re} \left\{ 6k d_0 \int_0^L F u_x (\overline{S}_x - \overline{U}) dx \right\}}_{\approx \mathcal{O}(|\lambda|^{-\ell/2})}. \quad (5.56)$$

Replacing (5.55) and (5.56) into (5.54) and using the fact $\ell = 2$ we obtain

$$3k \int_0^L F (|\mathcal{M}|^2)_x dx \lesssim \mathcal{O}(1). \quad (5.57)$$

Adding (5.52), (5.53) and (5.57) together with the fact that $\ell = 2$ we obtain the Claim 5.5.

□

Let $q_1, q_2 \in C^1([0, L])$ such that $0 \leq q_1(x) \leq 1$ and $0 \leq q_2(x) \leq 1$, for all $x \in [0, L]$ and

$$q_1(x) = \begin{cases} 1, & \text{if } x \in [0, \gamma_1], \\ 0, & \text{if } x \in [\gamma_2, L], \end{cases} \quad (5.58)$$

and

$$q_2(x) = \begin{cases} 0, & \text{if } x \in [0, \gamma_1], \\ 1, & \text{if } x \in [\gamma_2, L], \end{cases} \quad (5.59)$$

where $0 < \alpha < \gamma_1 < \gamma_2 < \beta < L$.

We consider

$$F = x q_1(x) + (x - L) q_2(x) = \begin{cases} x, & \text{if } x \in [0, \gamma_1], \\ (x - L), & \text{if } x \in [\gamma_2, L]. \end{cases} \quad (5.60)$$

Claim 5.6. We take $\ell = 2$. Let $F \in C^1([0, L])$ such that $F(0) = F(L) = 0$. The solution $W = (u, U, z, Z, s, S)^T \in \mathcal{D}(\mathcal{A})$ of (5.4) satisfies the following estimates

$$\Upsilon(\alpha + 4\varepsilon, \beta - 4\varepsilon) \approx \mathcal{O}(1) \quad \text{if } \ell = 2, \quad (5.61)$$

where

$$\begin{aligned} \Upsilon(\gamma_1, \gamma_2) &= \int_0^{\gamma_1} (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2) dx \\ &\quad + \int_{\gamma_2}^L (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2) dx \end{aligned} \quad (5.62)$$

for all $0 < \alpha < \gamma_1 < \gamma_2 < \beta < L$.

Proof. Taking $F = x q_1(x) + (x - L) q_2(x)$ in Claim 5.5 and using the definition of $d(x)$ with the fact that $0 < \alpha < \gamma_1 < \gamma_2 < \beta < L$ we obtain

$$\begin{aligned} &\int_0^{\gamma_1} (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2) dx \\ &\quad + \int_{\gamma_2}^L (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2) dx \\ &= - \int_{\gamma_1}^{\gamma_2} (q_1 + x q_{1x}) (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2 + d_0 |S_x - U|^2) dx \\ &\quad - \int_{\gamma_1}^{\gamma_2} (q_2 + (x - L) q_{2x}) (\rho_1 |\lambda u|^2 + k |u_x|^2 + \rho_2 |\lambda z|^2 + b |z_x|^2 + \rho_2 |\lambda s|^2 \\ &\quad + d_0 |S_x - U|^2) dx + d_0 \int_{\alpha}^{\gamma_2} q_1 |S_x - U|^2 dx + d_0 \int_{\gamma_1}^{\beta} q_2 |S_x - U|^2 dx. \end{aligned} \quad (5.63)$$

Taking $\gamma_1 = \alpha + 4\varepsilon$ and $\gamma_2 = \beta - 4\varepsilon$ into (5.63), using Claim 5.1-5.4 to $\ell = 2$ we obtain (5.61). □

Now we are in position to present the result of polynomial stability.

Theorem 5.7. *There is a constant $C > 0$ such that for every $U_0 \in \mathcal{D}(\mathcal{A})$, we have*

$$E(t) \leq \frac{C}{t} \|W_0\|_{\mathcal{D}(\mathcal{A})}, \quad t > 0. \quad (5.64)$$

Proof. From Claim 5.1-5.4 for $\ell = 2$ we have

$$\begin{cases} \int_{\alpha}^{\beta} |s_x|^2 dx \approx \mathcal{O}(\lambda^{-2}) \approx \mathcal{O}(1), & \int_{\alpha+\varepsilon}^{\beta-\varepsilon} |S|^2 dx \approx \mathcal{O}(1), & \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} |u_x|^2 dx \approx \mathcal{O}(1), \\ \int_{\alpha+2\varepsilon}^{\beta-2\varepsilon} |\lambda u|^2 dx \approx \mathcal{O}(1), & \int_{\alpha+3\varepsilon}^{\beta-3\varepsilon} |z_x|^2 dx \approx \mathcal{O}(1), & \int_{\alpha+4\varepsilon}^{\beta-4\varepsilon} |\lambda z|^2 dx \approx \mathcal{O}(1). \end{cases}$$

Thus, from (5.61) the above expression and the fact that $0 < \varepsilon < \frac{\beta-\alpha}{10}$ we conclude that $\|W\|_{\mathcal{H}} \approx \mathcal{O}(1)$ in $L^2(0, L)$, which we have a contradiction with (5.1). This implies that

$$\sup_{\lambda \in \mathbb{R}} \left\| (i\lambda I - \mathcal{A})^{-1} \right\|_{\mathcal{L}(\mathcal{H})} = \mathcal{O}(\lambda^2). \quad (5.65)$$

Using Borichev and Tomilov [5] the theorem follows. \square

Conclusion: We have studied the stabilization for a laminated beam with interfacial slip with one discontinuous local internal viscoelastic damping of Kelvin-Voigt type acting on the axial force under certain boundary conditions. We prove the strong stability of the system by using Arendt-Batty criteria. We prove that the total energy of the system decays polynomially.

Acknowledgments: The first author was partially financed by project Fondecyt 1191137.

REFERENCES

- [1] F. Alabau-Boussouira, Asymptotic behavior for Timoshenko beams subject to a single nonlinear feedback control, *Nonlinear Differ. Equations Appl.* 14, 643-669, 2007.
- [2] T. A. Apalara, On the stability of a thermoelastic laminated beam, *Acta. Math. Sci.* 39, 1517-1524, 2019.
- [3] W. Arendt and C. J. K. Batty, Tauberian theorems and stability of one-parameter semigroups, *Trans. Amer. Math. Soc.* 306, 837-852, 1998.
- [4] F. A. Amar-Khodja, J. E. Muñoz Rivera and R. Racke, Energy decay for Timoshenko systems of memory type, *J. Differ. Equations* 194, 82-115, 2003.
- [5] A. Borichev and Y. Tomilov, Optimal polynomial decay of functions and operator semigroups, *Math. Ann.* 347, 455-478, 2010.
- [6] X. G. Cao, D. Y. Liu and G. Q. Xu, Easy test for stability of laminated beams with structural damping and boundary feedback controls, *J. Dynamical Control Syst.* 13, 313-336, 2007.
- [7] B. Feng, D. S. Almeida Júnior and A. J. A. Ramos, Asymptotic stability for a laminated beam with structural damping and infinite memory, *Math. Nachr.* 294, 559-579, 2021.
- [8] B. Feng, T. F. Ma, R. N. Monteiro and C. A. Raposo, Dynamics of laminated timoshenko beams, *J. Dyn. Diff. Equat.* 30, 1489-1507, 2018.
- [9] D. X. Feng, D. H. Shi and W. Zhang, Boundary feedback stabilization of Timoshenko beam with boundary dissipation, *Sci. China. Ser. A: Math.* 41, 483-490, 1998.
- [10] A. Guesmia, Some well-posedness and general stability results in Timoshenko systems with infinite memory and distributed time delay, *J. Math. Phys.* 55, 2014. Article ID 081503
- [11] A. Guesmia and S. A. Messaoudi, General energy decay estimates of Timoshenko systems with frictional versus viscoelastic damping, *Math. Methods Appl. Sci.* 32, 2102-2122, 2009.
- [12] S. W. Hansen, A model for a two-layered plate with interfacial slip, In: *Control and Estimation of Distributed Parameter Systems: Nonlinear Phenomena* (Vorau, 1993), 143170. Birkhauser, Basel, 1994.

- [13] S. W. Hansen and R. Spies, Structural damping in a laminated beams due to interfacial slip, *J. Sound Vib.* 204, 183-202, 1997.
- [14] T. Kato, *Perturbation Theory for Linear Operators*, Springer, Berlin, Heidelberg, 1995.
- [15] A. Lo and N. E. Tatar, Stabilization of laminated beams with interfacial slip, *Electronic J. Differ. Equations* 129, 1-14, 2015.
- [16] A. Lo and N. E. Tatar, Exponential stabilization of a structure with interfacial slip, *Discrete Contin. Dyn. Syst.* 36, 6285-6306, 2016.
- [17] W. Liu, X. Kong and G. Li, Asymptotic stability for a laminated beam with structural damping and infinite memory, *Math. Mech. Solids* 25, 1979-2004, 2020.
- [18] Z. Liu and S. Zheng, *Semigroups Associated with Dissipative Systems*, Chapman & Hall, Vol. 398, London, 1999.
- [19] M. I. Mustafa, Boundary control of laminated beams with interfacial slip, *J. Math. Phys.* 59, 2018. Article ID 051508
- [20] C. Nonato, C. A. Raposo and B. Feng, Exponential stability for a thermoelastic laminated beam with nonlinear weights and time-varying delay, *Asympt. Anal.* Pre-press, 1-29, 2021.
- [21] A. Pazy, *Semigroups of Linear Operators and Applications to Partial Differential Equations*, Springer-Verlag, New York, 1983.
- [22] C. A. Raposo, Exponential stability for a structure with interfacial slip and frictional damping, *Appl. Math. Lett.* 53, 85-91, 2016.
- [23] C. A. Raposo, C. Nonato, O. V. Villagran and J. A. D. Chuquipoma, Global solution and exponential stability for a laminated beam with Fourier thermal law, *J. Partial Differ. Equ.* 33, 142-157, 2020.
- [24] C. A. Raposo, O. V. Villagrán, J. E. Muñoz Rivera and M. S. Alves, Hybrid laminated Timoshenko beam, *J. Math. Phys.* 58, (2017). Article ID 101512
- [25] S. P. Timoshenko, On the correction for shear of the differential equation for transverse vibrations of prismatic bars, *Philos. Mag. Ser. 6*, 744-746, 1921.
- [26] S. P. Timoshenko and J. M. Gere, *Mechanics of Materials*, D. Van Nostrand Company, Inc., New York, 1972.
- [27] J. M. Wang, G. Q. Xu and S. P. Yung, Exponential stabilization of laminated beams with structural damping and boundary feedback controls, *SIAM J. Control Optim.* 44, 1575-597, 2005.

OCTAVIO P. VERA VILLAGRÁN

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF BÍO-BÍO , CONCEPCIÓN, CHILE

E-mail address: `overa@ubiobio.cl`

CARLOS A. RAPOSO

DEPARTMENT OF MATHEMATICS, FEDERAL UNIVERSITY OF SÃO JOÃO DEL-REI, MINAS GERAIS, BRAZIL

E-mail address: `raposo@ufsj.edu.br`